

# Graphs, Linear Algebra, and Continuous Optimization

## Part III: Solving Exact Max Flow

**Aleksander Mądry**



# Breaking the $\Omega(n^{3/2})$ barrier

**Undirected** graphs and **approx.** answers ( $\Omega(n^{3/2})$  barrier still holds here)

[M '10]: **Crude approx. of** max flow **value** in **close to linear** time

[CKMST '11]: **(1- $\epsilon$ )-approx.** to max flow in  $\tilde{O}(n^{4/3}\epsilon^{-3})$  time

[LSR '13, S '13, KLOS '14, P '14]: **(1- $\epsilon$ )-approx.** in  $\tilde{O}(n\epsilon^{-2})$  time

**But:** What about the **directed** and **exact** setting?

[M '13]: Exact  $\tilde{O}(n^{10/7}) = \tilde{O}(n^{1.43})$ -time alg.

**Today** 

( $n$  = # of vertices,  $\tilde{O}()$  hides polylog factors)

# From electrical flows to **exact directed** max flow

**From now on:** All capacities are **1**,  $m=O(n)$   
and the value  $F^*$  of max flow is known

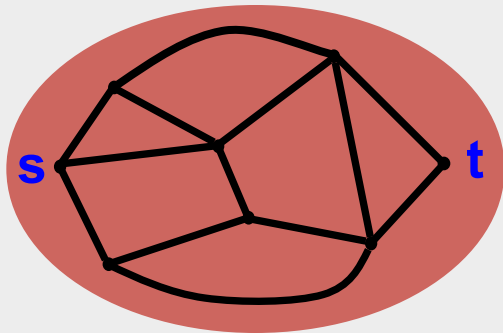
Why the progress on **approx. undirected** max flow does not apply to the **exact directed** case?

**Tempting answer:** Directed graphs are just different (for one, electrical flow is an undirected notion)

**But: exact directed** max flow reduces to **exact**

We need a more powerful intermediary

So, it is all about getting



**Key obstacle:** Gradient descent methods (like MWU) are inherently unable to deliver good enough accuracy

# (Path-following) Interior-point method (IPM)

[Dikin '67, Karmarkar '84, Renegar '88,...]

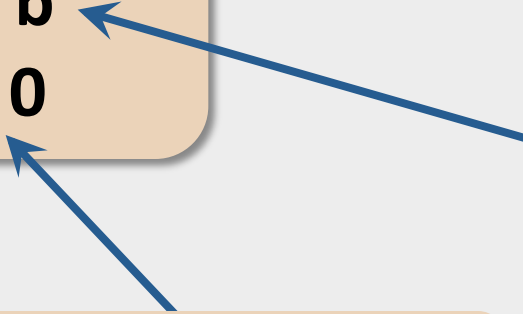
A powerful framework for solving general LPs (and more)

LP:  $\min c^T x$   
s.t.  $Ax = b$   
 $x \geq 0$

**Idea:** Take care of “hard” constraints by adding a “barrier” to the objective

“easy” constraints  
(use projection)

“hard” constraints



# (Path-following) Interior-point method (IPM)

[Dikin '67, Karmarkar '84, Renegar '88,...]

A powerful framework for solving general LPs (and more)

$$\begin{aligned} \text{LP}(\mu): \quad & \min \mathbf{c}^T \mathbf{x} - \mu \sum_i \log x_i \\ \text{s.t.} \quad & \mathbf{Ax} = \mathbf{b} \\ & \mathbf{x} \geq \mathbf{0} \end{aligned}$$

**Idea:** Take care of “hard” constraints by adding a “barrier” to the objective

**Observe:** The barrier term enforces  $\mathbf{x} \geq \mathbf{0}$  implicitly

**Furthermore:** for large  $\mu$ ,  $\text{LP}(\mu)$  is easy to solve and

$$\text{LP}(\mu) \rightarrow \text{original LP, as } \mu \rightarrow 0^+$$

## Path-following routine:

- Start with (near-)optimal solution to  $\text{LP}(\mu)$  for large  $\mu > 0$
- Gradually reduce  $\mu$  while maintaining the (near-)optimal solution to current  $\text{LP}(\mu)$

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**Idea:** Take care of “hard” constraints by adding a “barrier” to the objective

**Observe:** The barrier term enforces  $\mathbf{x} > \mathbf{0}$  implicitly.

Based on **second-order approx.**

$$\begin{aligned} f(\mathbf{x}+\mathbf{y}) \approx & f(\mathbf{x}) + \mathbf{y}^T \nabla f(\mathbf{x}) + \mathbf{y}^T \mathbf{H}_f(\mathbf{x}) \mathbf{y} \\ & + \text{projection on } \ker(\mathbf{A}) \end{aligned}$$

**Path-following routine:**

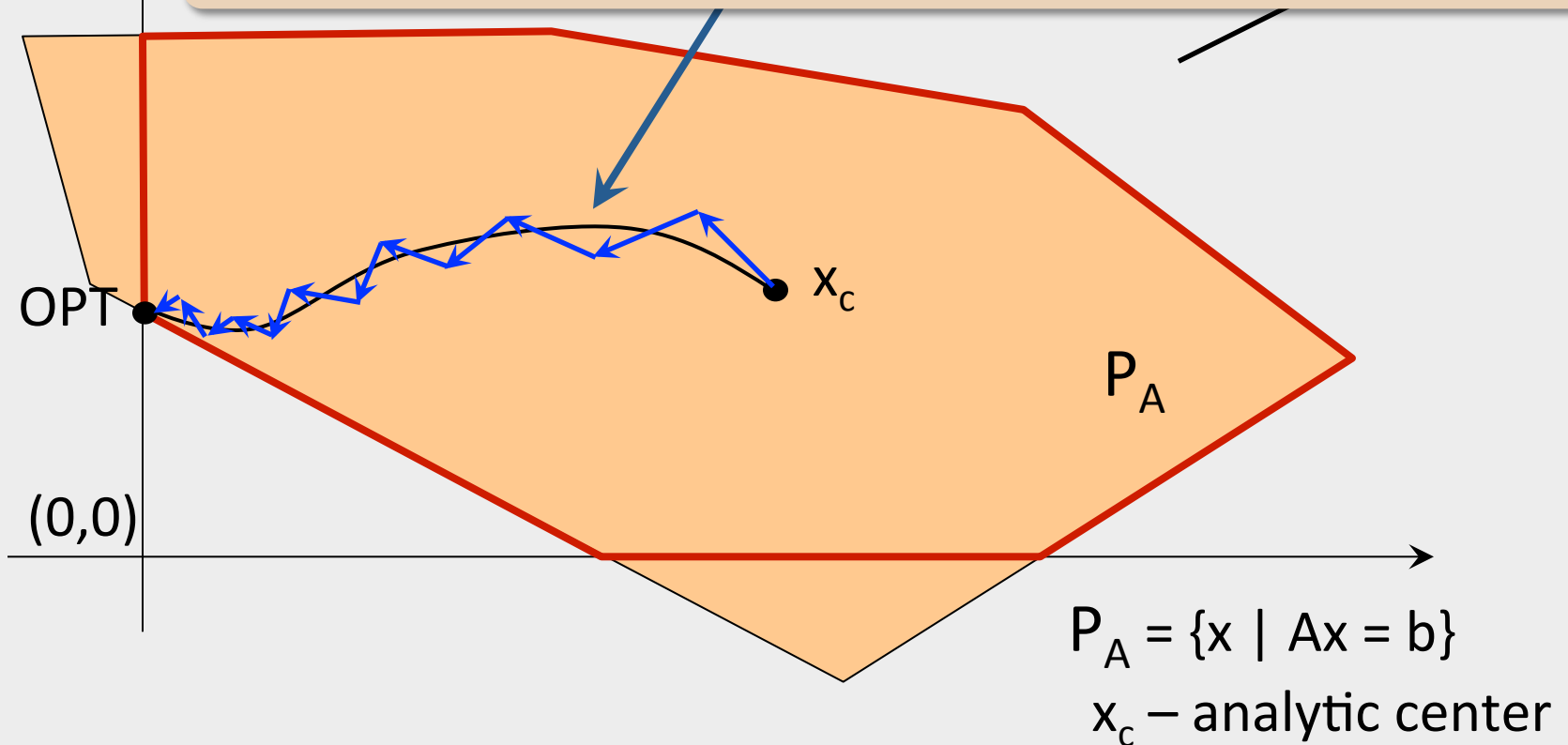
→ Maintain (near-)optimal solution

→ Repeat:

Set  $\mu' = (1-\delta)\mu$  and use **Newton's method** to compute from  $\mathbf{x}$  (near-)optimal solution to  $\text{LP}(\mu')$

**Key point:** Choosing step size  $\delta$  sufficiently small ensures  $\mathbf{x}$  is close to optimum for  $\text{LP}(\mu')$  → Newton's method convergence **very** rapid

**central path** = optimal solutions to  $\text{LP}(\mu)$  for all  $\mu > 0$



### Path-following routine:

- Start with (near-)optimal solution to  $\text{LP}(\mu)$  for large  $\mu > 0$
- Gradually reduce  $\mu$  (via **Newton's method**) while maintaining the (near-)optimal solution to current  $\text{LP}(\mu)$



# Can we use IPM to get a faster max flow alg.?

**Conventional wisdom:** This will be too slow!

→ Each **Newton's step** = solving a linear system  $O(n^\omega) = O(n^{2.373})$  time  
(prohibitive!)

**But:** When solving **flow problems** – only  $\tilde{O}(m)$  time [DS '08]

**Fundamental question:** What is the number of iterations?

[Renegar '88]:  $O(m^{1/2} \log \varepsilon^{-1})$

**Unfortunately:** This gives only an  $\tilde{O}(m^{3/2})$ -time algorithm

**Improve the  $O(m^{1/2})$  bound?**

Although believed to be **very** suboptimal,  
its improvement is a major challenge

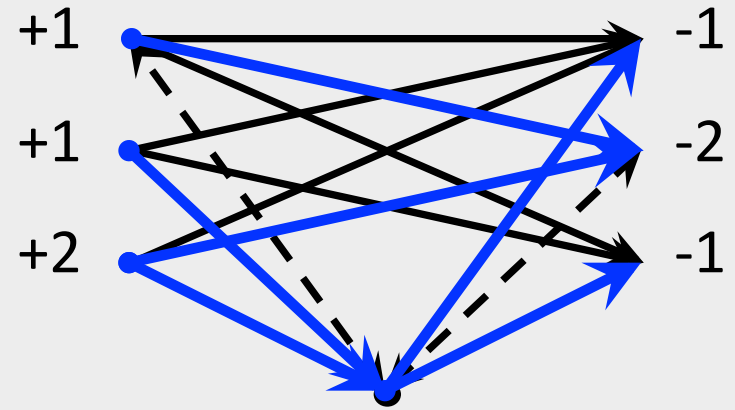
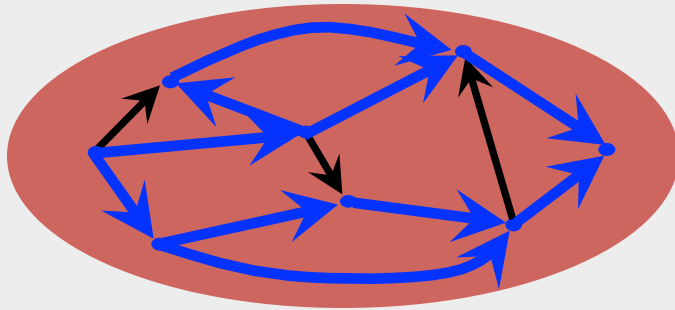


# The Max Flow algorithm

(Self-contained, but can be seen as a variation on IPM)

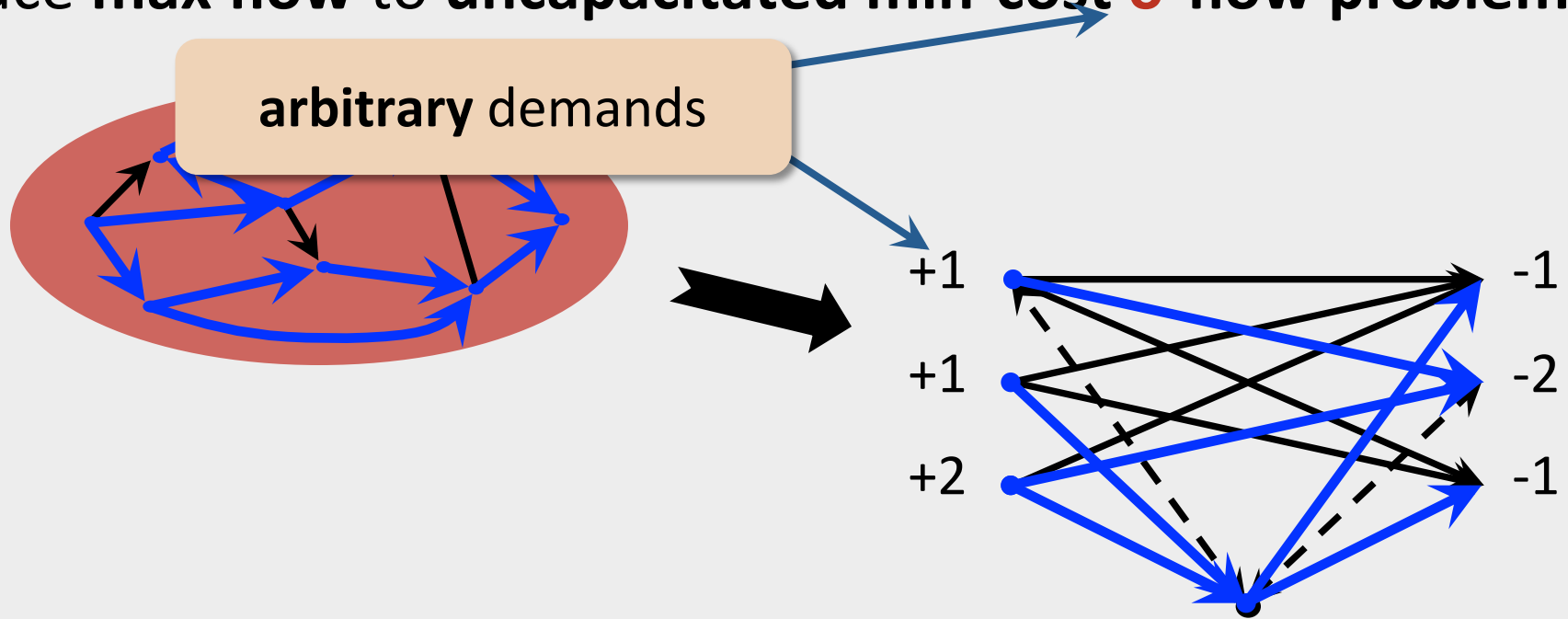
# From Max Flow to Min-cost Flow

Reduce max flow to uncapacitated min-cost  $\sigma$ -flow problem



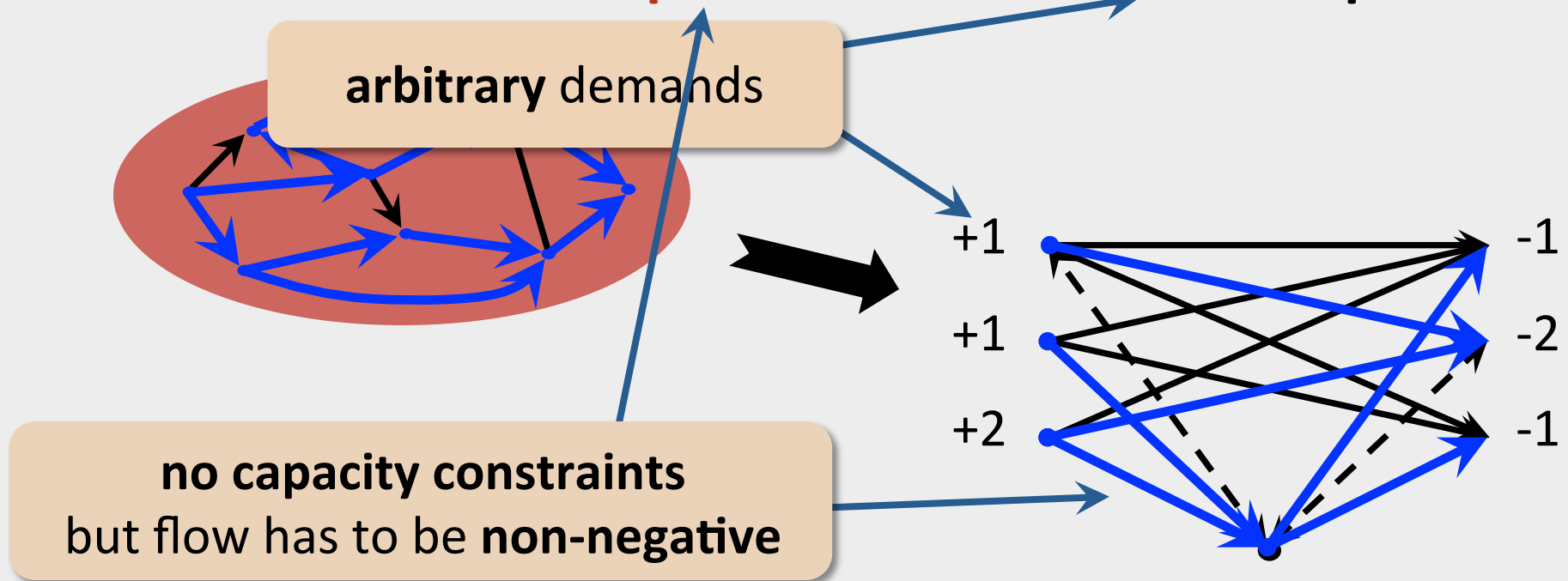
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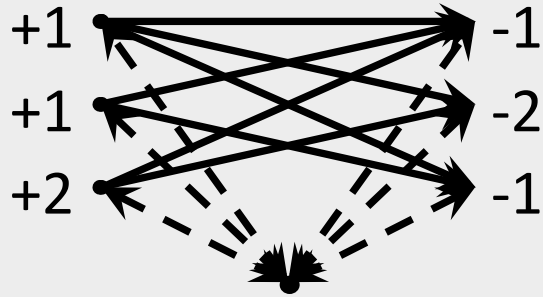
# From Max Flow to Min-cost Flow

Reduce max flow to **uncapacitated** min-cost  $\sigma$ -flow problem



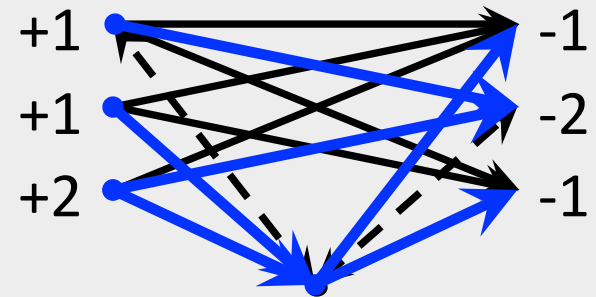
**Result:** Feasibility  $\rightarrow$  Optimization  
+ special structure

# Solving Min-Cost Max Flow Instance



Our approach is **primal-dual**

→ **Primal solution:  $\sigma$ -flow  $f$**   
(feasibility: all  $f_e$  are  $\geq 0$ )

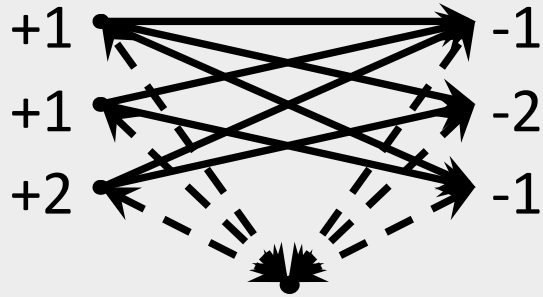


→ **Dual solution: embedding  $y$  into real line**  
(feasibility: all slacks  $s_e$  are  $\geq 0$ )

“No arc is too stretched”

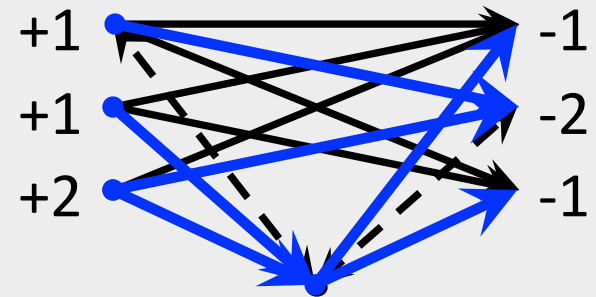


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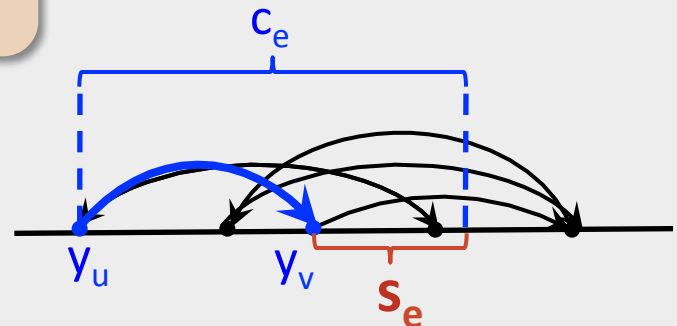
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“No arc is too stretched”



# Solving Min-Cost Max Flow Instance

## Our Goal:

Get  $(f, y)$  with small **duality gap**  $\sum_e f_e s_e$

**Our Approach:** Iteratively improve maintained solution while enforcing an **additional** constraint

## Centrality:

$f_e s_e \approx \mu$ , for all  $e$   
(with  $\mu$  being progressively smaller)

“Make all arcs have similar contribution to the duality gap”

(Maintaining **centrality** = following the **central path**)



# Taking an Improvement Step

So far, our approach is fairly standard

## **Crucial Question:**

How to improve the quality of maintained solution?

## **Key Ingredient:**

Use electrical flows

# Taking an Improvement Step

Let  $(\mathbf{f}, \mathbf{y})$  be a (centered) primal-dual solution

**Key step:** Compute **electrical  $\sigma$ -flow  $\mathbf{f}^+$**  with  $r_e := s_e / f_e$

**Primal improvement:** Set  $\mathbf{f}' := (1-\delta)\mathbf{f} + \delta\mathbf{f}^+$

**Dual improvement:** Use **voltages  $\boldsymbol{\varphi}$**  inducing  $\mathbf{f}^+$  (via Ohm's Law)  
Set  $\mathbf{y}' := \mathbf{y} + \delta(1-\delta)^{-1} \boldsymbol{\varphi}$

**Can show:** When terms **quadratic** in  $\delta$  are ignored

$$f'_e s'_e \approx (1-\delta) \mu = \mu' \\ \text{for each } e$$



(i.e., **duality gap** decreases by  $(1-\delta)$  and **centrality** is preserved)

How big  $\delta$  can we take to have this approx. hold?

# Lowerbounding $\delta$

Can show:

$\delta^{-1}$  is bounded by  $O(|\rho|_4)$   
where  $\rho_e := |f_e^+|/f_e$

$|\rho|_4$  measures  
how different  $f^+$  and  $f$  are

How to bound  $|\rho|_4$ ?

Idea: Bound  $|\rho|_2 \geq |\rho|_4$  instead

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$(|\rho|_2 \geq |\rho|_4)$

Centrality: Tying  $|\rho|_2$  to  $E(f^+)$

$$f_e s_e \approx \mu \rightarrow r_e = s_e/f_e \approx \mu/(f_e)^2$$



$$E(f^+) \approx \mu (|\rho|_2)^2$$

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$$f_e s_e \approx \mu \rightarrow r_e = s_e/f_e \approx \mu/(f_e)^2$$



$$E(f^+) = \sum_e r_e (f_e^+)^2 \approx \sum_e \mu (f_e^+/f_e)^2 = \mu \sum_e (\rho_e)^2 = \mu (|\rho|_2)^2$$

So, we can focus on bounding  $E(f^+)$

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How to bound  $|\rho|_2$ ?

$$(|\rho|_2 \geq |\rho|_4)$$

How to bound  $E(f^+)$ ?

$$(E(f^+) \approx \mu (|\rho|_2)^2)$$

**Idea:** Use energy-bounding argument  
we used in the undirected case

**Claim:**  $E(f^+) \leq \mu m$

**Proof:** Note that  $E(f) = \sum_e r_e (f_e)^2 \approx \sum_e \mu (f_e/f_e)^2$

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How to bound  $|\rho|_2$ ? ( $|\rho|_2 \geq |\rho|_4$ )

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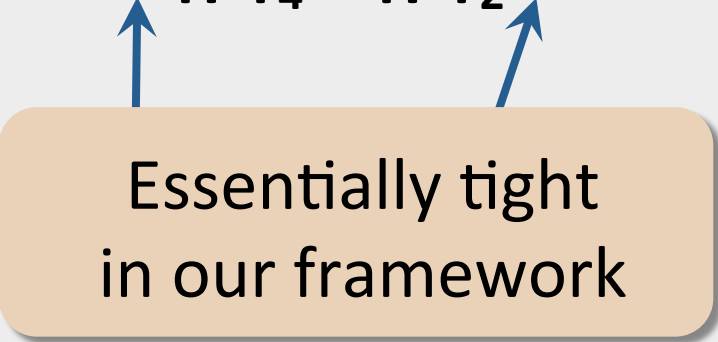
**Result:** Bounding  $\delta^{-1} \leq |\rho|_4 \leq |\rho|_2 \leq (E(f^+)/\mu)^{1/2} \leq m^{1/2}$

$E(f^+) \leq E(f) \approx \mu m$

This recovers the canonical  $O(m^{1/2})$ -iterations bound  
for **general IPMs** and gives the  $\tilde{O}(m^{3/2} \log U)$  algorithm

# Going beyond $\Omega(m^{1/2})$ barrier

Our reasoning before:  $\delta^{-1} \leq |\rho|_4 \leq |\rho|_2 \leq m^{1/2}$



Essentially tight  
in our framework



# Going beyond $\Omega(m^{1/2})$ barrier

Our reasoning before:  $\delta^{-1} \leq |\rho|_4 \leq |\rho|_2 \leq m^{1/2}$

When does  $|\rho|_4 \approx |\rho|_2$ ?



This part we need  
to improve

# Going beyond $\Omega(m^{1/2})$ barrier

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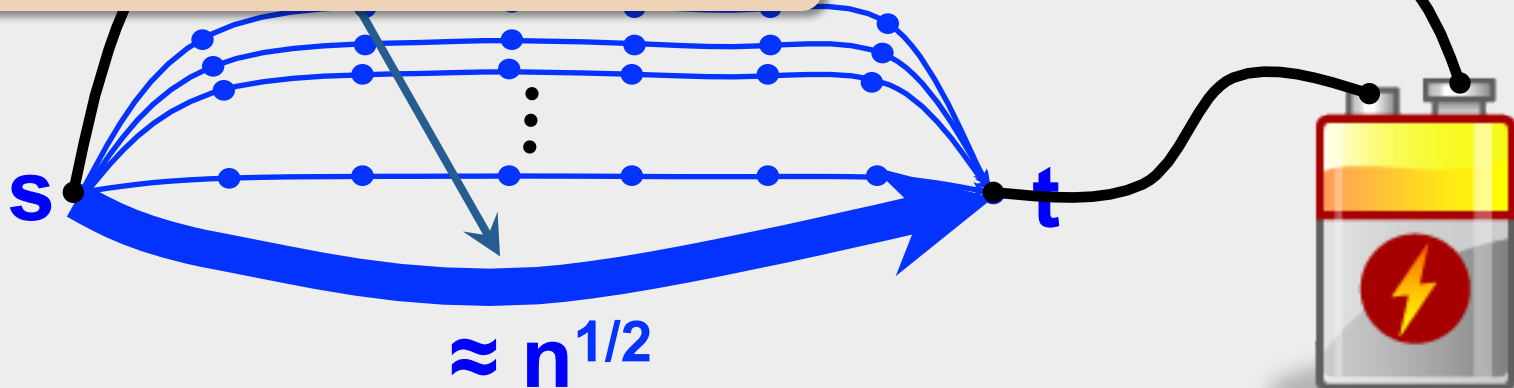
When does  $|\rho|_4 \approx |\rho|_2$ ? **Answer:** If most of the norm of  $\rho$  is focused on only a few coordinates

**Translated to our setting:**  $|\rho|_4 \approx |\rho|_2$  if most of the energy of  $f^+$  is contributed by only a few arcs

Can this happen?

Unfortunately, yes

Contributes most of the energy



# Going beyond $\Omega(m^{1/2})$ barrier

Our reasoning before:  $\delta^{-1} \leq \|\rho\|_4 \leq \|\rho\|_2 \leq m^{1/2}$

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This is the **only** part where **unit-capacity** assumption is needed

(in principle, tight)

open too often

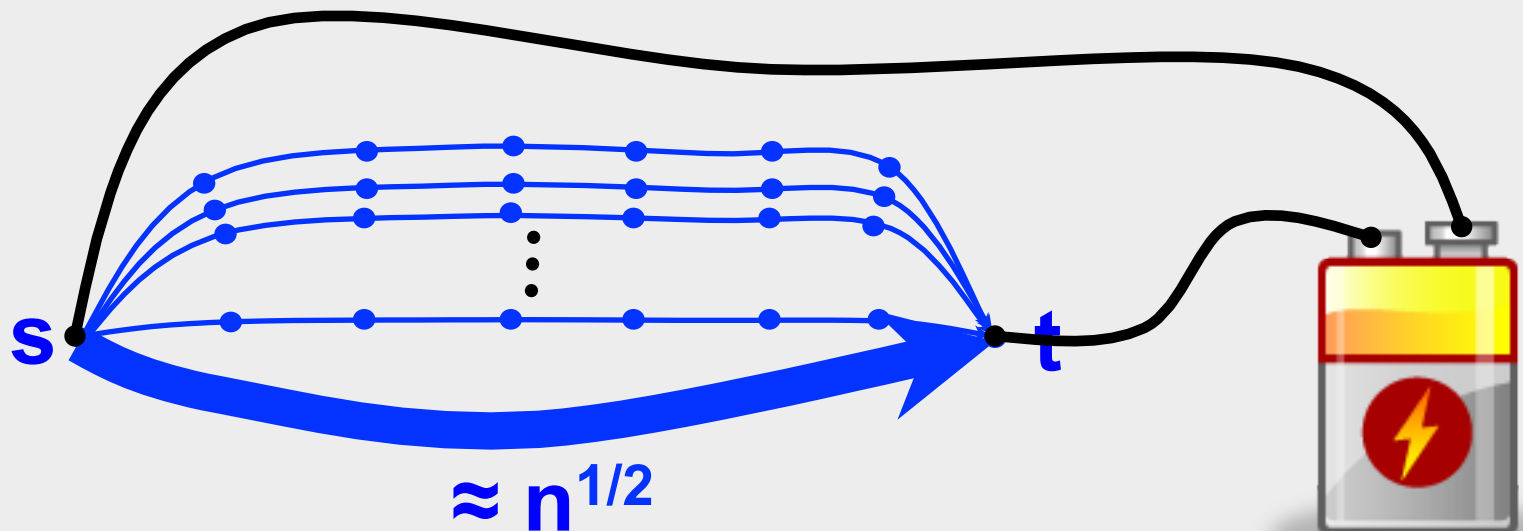
**Method: Very careful perturbation** of the solution  
+ certain **preconditioning**

# Going beyond $\Omega(m^{1/2})$ barrier

**Problematic case:** When most of the energy of  $f^+$  is contributed by only a few arcs

How can we ensure that this is not the case?

We already faced such problems in the undirected setting!



# Going beyond $\Omega(m^{1/2})$ barrier

**Problematic case:** When most of the energy of  $f^+$  is contributed by only a few arcs

How can we ensure that this is not the case?

We already faced such problems in the undirected setting!

**Our approach then:** Keep removing high-energy edges

**To show this works:** Use the energy of the electrical flow as a potential function

- Energy **can only increase** and obeys global upper bound
- Each time removal happens  $\rightarrow$  energy **increases by a lot**

**Problems:** In our framework, arc removal is **too drastic** and the energy of  $f^+$  is **highly non-monotone**

# Going beyond $\Omega(m^{1/2})$ barrier

How to deal with these problems?

→ Enforce a **stronger** condition than just that  $|\rho|_4$  is small (“smoothness”: restrict energy contributions of arc subsets)

**Key fact:  $f^+$  smooth** → energy does **not** change too much (so, it is a good potential function again)

→ To enforce this, keep **stretching** the offending arcs (stretch = increase length by  $s_e$  - this doubles the resistance  $r_e = s_e/f_e$ )

As long as  $s_e$  is small for stretched arcs, the resulting perturbation of lengths can be corrected at the end

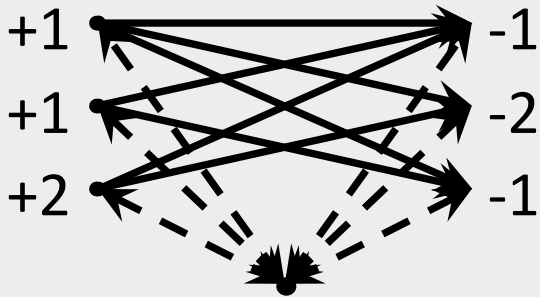
**Remaining question:** How to handle arcs with **large  $s_e$** ?

# Going beyond $\Omega(m^{1/2})$ barrier

**Observation:** As  $f_e s_e \approx \mu$ , large  $s_e \rightarrow$  small flow  $f_e$   
and thus  $r_e = s_e / f_e \approx \mu / f_e^2$  is pretty large

$\rightarrow$  **For such arcs:** contributing a lot of energy implies  
high effective resistance

**Idea:** Precondition  $(f, y)$  so as no arc has too high effect. resist.

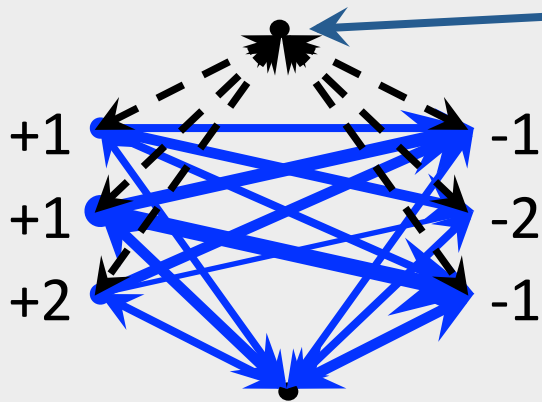


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Auxiliary star graph

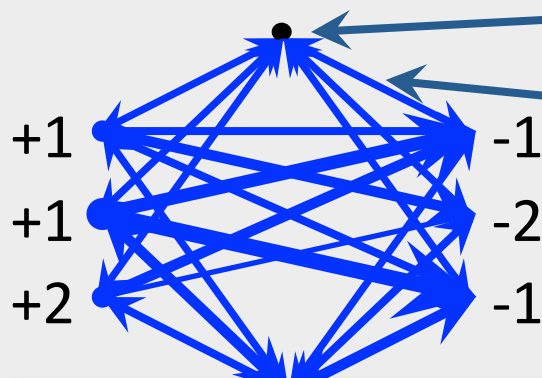


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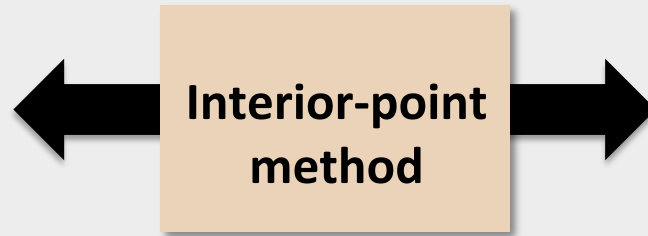
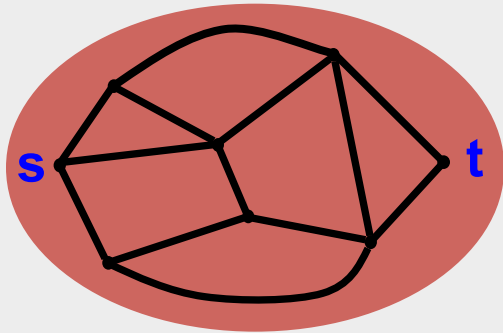
Trivial circulations on each pair of arcs

**Can show:** After doing that, no arc with large  $s_e$

**Putting these two techniques together + some work:**  
 $\tilde{O}(m^{3/7})$ -iterations convergence follows

# **Conclusions and the Bigger Picture**

# Maximum Flows and Electrical Flows



**Elect. flows + IPMs** → A powerful new approach to **max flow**

Can this lead to a **nearly-linear time** algorithm for the **exact directed** max flow?

We seem to have the “critical mass” of ideas



**Elect. flows** = next generation of “spectral” tools?

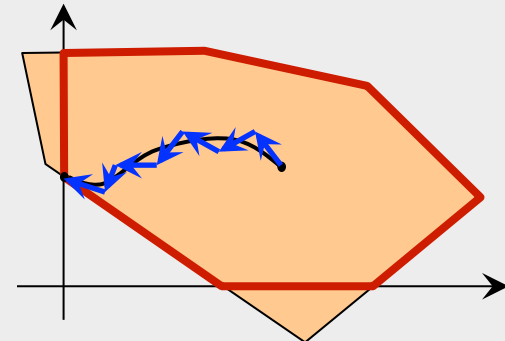
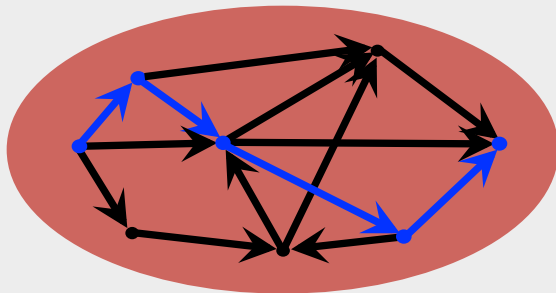
- Better “spectral” graph partitioning,
- Algorithmic grasp of random walks,
- ...

**Grand challenge:** Can we make algorithmic graph theory run in nearly-linear time?

**New “recipe”:** Fast alg. for **combinatorial** problems via **linear-algebraic** tools + **continuous opt.** methods

How about applying this framework to other graph problems that “got stuck” at  $O(n^{3/2})$ ?  
(min-cost flow, general matchings, negative-lengths shortest path...)

**Second-order/IPM-like methods:**  
the next frontier for fast (graph) algorithms?



# Max Flow and Interior-Point Methods

**Contributing back:** Max flow and electrical flows as a lens for analyzing general IPMs?

Our techniques can be lifted to the general LP setting

We can solve **any** LP within  $\tilde{O}(m^{3/7}L)$  iterations  
**But:** this involves **perturbing** of this LP

Some (seemingly) new elements of our approach:

- Better grasp of  $\ell_2$  vs.  $\ell_4$  interplay wrt the step size  $\delta$
- Perturbing the central path when needed
- Usage of non-local convergence arguments

Can this lead to breaking the  $\Omega(m^{1/2})$  barrier for all LPs?

[Lee Sidford '14]:  $\tilde{O}(\text{rank}(\mathbf{A})^{1/2})$  iteration bound

# Bridging the Combinatorial and the Continuous

paths, trees, partitions,  
routings, matchings,  
data structures...



matrices, eigenvalues,  
linear systems, gradients,  
convex sets...

**Powerful approach:** Exploiting the interplay of the two worlds

Some other early “success stories” of this approach:

- Spectral graph theory aka the “eigenvalue connection”
- Fast SDD/Laplacian system solvers
- Graph sparsification, random spanning tree generation

...and this is just the beginning!

**Thank you**

**Questions?**